

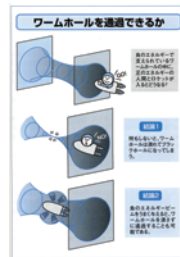
Wormhole Evolutions in Higher-dim. Gravity --- Effects of Gauss-Bonnet gravity terms



Hisaaki Shinkai & Takashi Torii
(Osaka Inst. Technology, Japan)
真貝寿明 & 鳥居隆 (大阪工業大学)

Outline & Summary

- (a) "Fate of Morris-Thorne (Ellis) wormhole" was numerically investigated in 2002. [HS & Hayward, PRD66, 044005]. The fate is either black-hole collapse or inflationary expansion, depending on the exceeded energy.
- (b) The higher-dimensional Ellis wormhole solutions are obtained. Perturbation study suggests instability. Numerical evolutions confirm its instability.
- (c) The wormholes in 5-dim. Gauss-Bonnet gravity are numerically obtained. Evolutions suggest that positive GB term accelerates black hole collapse.



Motivations

Why wormholes?

- They make great science fiction – short cuts between otherwise distant regions. Morris & Thorne 1988, Sagan "Contact" etc
- They increase our understanding of gravity when the usual energy conditions are not satisfied, due to quantum effects (Casimir effect, Hawking radiation) or alternative gravity theories, brane-world models etc.
- They are very similar to black holes – both contain (marginally) trapped surfaces and can be defined by trapping horizons (TH).

Wormhole ≡ Hypersurface foliated by marginally trapped surfaces

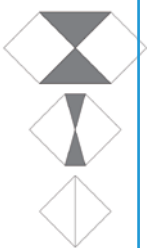
- BH and WH are interconvertible?

- How the stability changes in 5-d GR?
- How the stability changes in Gauss-Bonnet gravity?

BH and WH are interconvertible? (New Duality?)

S.A. Hayward, Int. J. Mod. Phys. D 8 (1999) 373

- They are very similar – both contain (marginally) trapped surfaces and can be defined by trapping horizons (TH)
- Only the causal nature of the THs differs, whether THs evolve in plus / minus density.



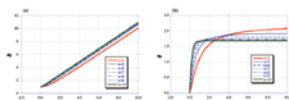
	Black Hole	Wormhole
Locally defined by	Achronal (spatial/null) outer TH ⇒ 1-way traversable	Temporal (timelike) outer THs ⇒ 2-way traversable
Einstein eqs.	Positive energy density normal matter (or vacuum)	Negative energy density "exotic" matter
Appearance	occur naturally	Unlikely to occur naturally, but constructible ???

N-dim. Ellis Wormhole sol.

in prep.

A Wormhole Solution (n-Dim, massless ghost scalar)

- spherical symmetry.
- $ds^2 = -f(t,r)e^{-2\alpha(t,r)}dt^2 + f(t,r)^{-1}dr^2 + R(t,r)^2d\Omega^2$ (1)
- with massless ghost scalar field ϕ .
- static, $f \equiv 1$, and throat radius $R(0) = a$; Just solve
- $\frac{d^2R}{dr^2} = \frac{(n-3)a^{2(n-3)}}{R^{2n-5}}$ (2)
- $\frac{d\phi}{dr} = \sqrt{(n-2)(n-3)}\frac{a^{n-3}}{R^{n-2}}$ (3)



The n-dimensional wormhole solutions. The circumference radius R (a) and the scalar field phi (b) are plotted as a function of radial coordinate r. The cases of n = 4-10 are shown.

Perturbation Analysis

$$f(t,r) = f_0(r) + \epsilon f_1(r)e^{i\omega t},$$

$$\delta(t,r) = \delta_0(r) + \epsilon \delta_1(r)e^{i\omega t},$$

$$R(t,r) = R_0(r) + \epsilon R_1(r)e^{i\omega t},$$

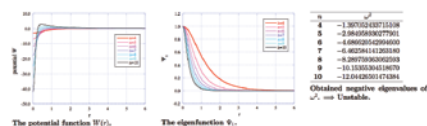
$$\phi(t,r) = \phi_0(r) + \epsilon \phi_1(r)e^{i\omega t},$$

- Gauge invariant quantity $\psi_1 = R_0^{n/2}(\phi_1 - \frac{\delta_1}{R_0})$
- regularized variable $\Psi_1 = D_r \psi_1$ where $D_r = \frac{d}{dr} - \frac{\psi_1'}{\psi_1}$, the master equation becomes

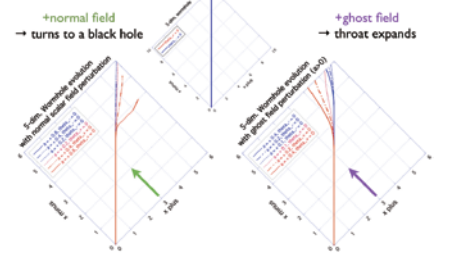
$$-\Psi_1'' + W(r)\Psi_1 = \omega^2\Psi_1, \quad (8)$$

where

$$W(r) = -\frac{1}{4R_0^3} \frac{\beta(n-2)^2}{R_0^{2(n-3)}} - (n-4)(n-6). \quad (9)$$



Numerical Evolutions



Results in 4-dim. GR

PRD66 (2002) 044005

Bifurcation of the horizons

-- go to a Black Hole or inflationary expansion

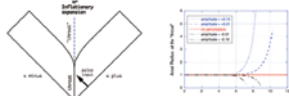


Figure 6: Phase diagram of the horizon...

Ghost pulse input - Bifurcation of the horizons

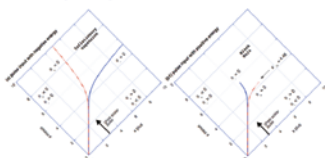


Figure 8: Bifurcation diagrams...

Travel through a Wormhole - with Maintenance Operations!

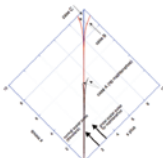
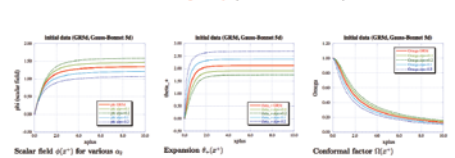


Figure 10: A field of wormhole solutions...

WH evolution in 5-dim. GB

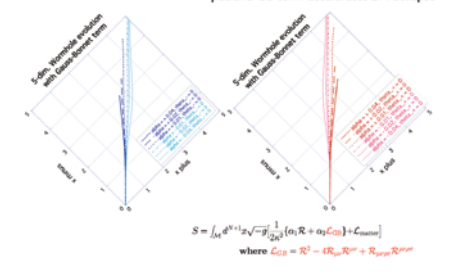
in prep.

Wormholes in Gauss-Bonnet gravity (initial data on x^0)



WH evolution in 5D Gauss-Bonnet gravity

positive GB term accelerates BH collapse



Field Eqs.

Gauss-Bonnet gravity

$$S = \int d^4x \sqrt{-g} \left[\frac{1}{2\kappa^2} (R + \alpha_1 C_{\mu\nu}C^{\mu\nu} + \alpha_2 C_{\mu\nu\rho\sigma}C^{\mu\nu\rho\sigma}) + C_{\text{matter}} \right] \quad (1)$$

$$\text{where } C_{GB} = R^2 - 4R_{\mu\nu}R^{\mu\nu} + R_{\mu\nu\rho\sigma}R^{\mu\nu\rho\sigma}$$

$$\alpha_1 G_{\mu\nu} + \alpha_2 H_{\mu\nu} + g_{\mu\nu} \Lambda = \kappa^2 T_{\mu\nu} \quad (2)$$

$$\text{where } H_{\mu\nu} = 2[R_{\mu\nu} - 2R_{\mu\nu}R^{\nu\sigma} - 2R^{\sigma\lambda}R_{\mu\lambda} + R_{\rho\sigma\lambda\tau}R^{\rho\sigma\lambda\tau}] - \frac{1}{3}g_{\mu\nu}C_{GB}$$

- Features
 - has GR correction terms from String Theory.
 - has two solution branches (GR/non-GR).
 - is expected to have singularity avoidance feature. (but has never been demonstrated.)

- new topic in numerical relativity. (S Golec & T Piran, PRD 85 (2012) 104015; F Izaurieta & E Rodriguez, 1207.1496; N Deppe+ 1208.8250)

Assumptions

- 5-dim. Spherical Symmetry
- Dual-null coordinate

$$ds^2 = -2e^{2\psi}dv^2 + dx^2 + r^2(e^{2\theta}d\Omega^2) \quad (3)$$

$$\text{conformal factor } \Omega = \frac{1}{r} \quad (4)$$

$$\text{expansions } \theta_{\pm} = 3\partial_r \psi - 3\Omega^{-1}\partial_v \Omega \quad (5)$$

$$\text{infinities } \nu_{\pm} = \partial_r f \quad (6)$$

$$\text{momenta of } \phi \quad \nu_{\pm} = r\partial_r \phi = \Omega^{-1}\partial_v \phi \quad (7)$$

$$\text{momenta of } \psi \quad \pi_{\pm} = r\partial_r \psi = \Omega^{-1}\partial_v \psi \quad (8)$$

- matter = normal field $\psi(u,v)$ and/or ghost field $\phi(u,v)$

$$T_{\mu\nu} = T_{\mu\nu}^{\psi} + T_{\mu\nu}^{\phi} = \left[\psi_{,\mu}\psi_{,\nu} - g_{\mu\nu} \left(\frac{1}{2}(\nabla\psi)^2 + V_1(\psi) \right) \right] + \left[\phi_{,\mu}\phi_{,\nu} - g_{\mu\nu} \left(\frac{1}{2}(\nabla\phi)^2 + V_2(\phi) \right) \right] \quad (9)$$

Field Equations in 5-Dim. Gauss-Bonnet gravity

$$\eta = \Omega^2 (e^{-f} + \frac{2}{3}\partial_r \psi), \quad \tilde{A} = (\alpha_1 + 4\alpha_2\nu_{\pm}), \quad B = \kappa^2 T_{\pm\pm} - \kappa^4 A$$

$$x^{\pm}\text{-direction } \partial_{\pm} \quad x^r\text{-direction } \partial_r$$

$$\partial_v \Omega = -\frac{1}{2}\Omega^2 \partial_v \psi \quad (10)$$

$$\partial_v \nu_{\pm} = -\nu_{\pm} \partial_v \nu_{\pm} - \frac{1}{3\Omega} \partial_v^2 \nu_{\pm} \quad (11)$$

$$\partial_v \nu_{\pm} = \frac{1}{\Omega} (\partial_v^2 \nu_{\pm} - 3\nu_{\pm} \partial_v \eta) \quad (12)$$

$$\partial_v f = \nu_{\pm} \quad (13)$$

$$\partial_v \nu_{\pm} = \frac{\nu_{\pm}}{\tilde{A}} \left[\eta - \frac{4(3\alpha_2 \eta - B)}{3\tilde{A}} \right] + \frac{\kappa^2 T_{\pm\pm} (\Omega^2 - \tilde{A})}{\kappa^2 \tilde{A}} + \frac{8\alpha_2}{3\tilde{A}^2} \left[\nu_{\pm}^2 (3\alpha_2 \eta - B) - \kappa^2 T_{\pm\pm} \nu_{\pm} \right] \quad (14)$$

$$\partial_v \psi = \Omega \pi_{\pm} \quad (15)$$

$$\partial_v \pi_{\pm} = \Omega \pi_{\pm} \quad (16)$$

$$\partial_v \pi_{\pm} = -\frac{1}{2} \partial_v^2 \pi_{\pm} - \frac{1}{3} \partial_v^3 \pi_{\pm} - \frac{1}{3\Omega} \frac{d\Omega}{dv} \partial_v \pi_{\pm} \quad (17)$$

$$\partial_v \pi_{\pm} = -\frac{1}{2} \partial_v^2 \pi_{\pm} - \frac{1}{3} \partial_v^3 \pi_{\pm} - \frac{1}{3\Omega} \frac{d\Omega}{dv} \partial_v \pi_{\pm} \quad (18)$$

Energy-momentum tensor

$$T_{\pm\pm} = \Omega^2 (\pi_{\pm}^2 - \pi_{\pm}^2) \quad (10)$$

$$T_{\pm r} = -\Omega^2 (\pi_{\pm}^2 - \pi_{\pm}^2) \quad (11)$$

$$T_{\pm v} = -\kappa^2 (V_1(\psi) + V_2(\phi)) \quad (12)$$

$$T_{\pm r} = e^f (\pi_{\pm} \nu_{\pm} - \pi_{\pm} \nu_{\pm}) - \frac{1}{3\Omega} (V_1(\psi) - V_2(\phi)) \quad (13)$$