

ワームホールの不安定性

真貝寿明 (大阪工業大学情報科学部) 鳥居隆 (工学部)
Hisaaki Shinkai & Takashi Torii (Osaka Inst. Technology)

Outline & Summary

タイムマシンの原理としても有名な Morris-Thorne のワームホール解 (Ellis WH 解) はゴーストスカラー場で構成される仮想的な一般相対論の解である。静的な仮定で導かれた解だが、著者らは動的に不安定であることを以前報告した【Shinkai & Hayward, Phys. Rev. D 66 (2002) 044005】。今回は、高次元に拡張したり、宇宙項を入れたとしても同様に不安定であることを報告する。

我々は、 n 次元に拡張した時空中で Ellis WH 解を求め、摂動に対して不安定であることを見つけた【Torii & Shinkai, PRD88 (2013) 064023】。数値計算により、不安定であることを確かめ、WH の喉 (throat) は、与えるゆらぎのエネルギーの正負によって、ブラックホール (BH) に転じるか、あるいはインフレーション的に膨張することがわかった。BH に転じる場合、与えるゆらぎは小さくても形成される BH 質量には最小値が存在し、次元が上がれば最小値は大きくなる。宇宙項の正負によっても WH は BH または拡大する。Gauss-Bonnet 重力理論では、BH 形成は抑えられるようだが、定常解ではない。総じて、このような単純なワームホールは不安定であると結論でき、観測するのは難しそうである。



Motivations

Why wormholes?

- They make great science fiction – short cuts between otherwise distant regions. Morris & Thorne 1988, Sagan "Contact" etc
- They increase our understanding of gravity when the usual energy conditions are not satisfied, due to quantum effects (Casimir effect, Hawking radiation) or alternative gravity theories, brane-world models etc.
- They are very similar to black holes – both contain (marginally) trapped surfaces and can be defined by trapping horizons (TH).

Wormhole \equiv Hypersurface foliated by marginally trapped surfaces

- BH and WH are interconvertible?
- New duality?

- How the stability changes in 5-d GR?
- How the stability changes in Gauss-Bonnet gravity?

BH and WH are interconvertible? (New Duality?)

S.A. Hayward, Int. J. Mod. Phys. D 8 (1999) 373

- They are very similar – both contain (marginally) trapped surfaces and can be defined by trapping horizons (TH)
- Only the causal nature of the THs differs, whether THs evolve in plus / minus direction.



	Black Hole	Wormhole
Locally defined by	Achronal (spatial/null) outer TH \Rightarrow 1-way traversable	Temporal (timelike) outer TH \Rightarrow 2-way traversable
Einstein eqs.	Positive energy density normal matter (or vacuum)	Negative energy density "exotic" matter
Appearance	occur naturally	Unlikely to occur naturally, but constructible ???

N-dim. Ellis Wormhole sol.

A Wormhole Solution (n-Dim, massless ghost scalar)

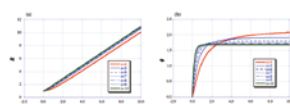
- spherical symmetry.

$$ds^2 = -f(t,r)c^{2\alpha(t,r)}dt^2 + f(t,r)^{-1}dr^2 + R(t,r)^2d\Omega^2 \quad (1)$$

- with massless ghost scalar field ϕ .
- static, $f \equiv 1$, and throat radius $R(0) = a$; Just solve

$$\frac{d^2R}{dr^2} = \frac{(n-3)2^{2(n-3)}}{R^{2(n-3)}} \quad (2)$$

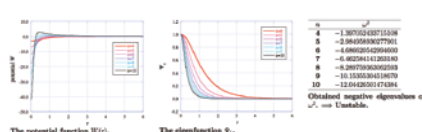
$$\frac{d\phi}{dr} = \sqrt{(n-2)(n-3)}\frac{1}{R^{n-2}} \quad (3)$$



The n-dimensional wormhole solutions. The circumferential radius $R(r)$ and the scalar field $\phi(r)$ are plotted as a function of radial coordinate r . The cases of $n=4-6$ are shown.

Torii & HS, PRD88 (2013), 064023

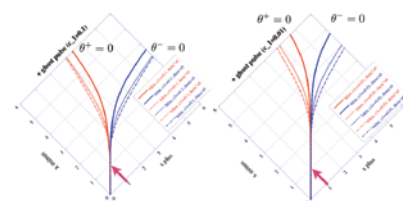
Perturbation Analysis



摂動に対して不安定なモードが必ず存在する

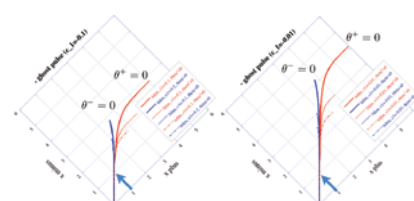
WH evolution in 4, 5, 6-dim. GR in prep.

4d 5d 6d GR ghost pulse (additional amp.) input



negative energy input \rightarrow throat inflates

4d 5d 6d GR ghost pulse (subtract amp.) input



positive energy input \rightarrow BH formation

Results in 4-dim. GR

PRD66 (2002) 044005

Bifurcation of the horizons – go to a Black Hole or inflationary expansion

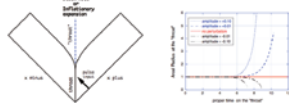


Figure 4: Partial Penrose diagram of the initial spacetime. Figure 5: Area radius r of the "throat" $r = r^*$ plotted as a function of proper time. Additional negative energy causes inflationary expansion, which reduces negative energy source collapse to a black hole and causes expansion.

Ghost pulse input – Bifurcation of the horizons

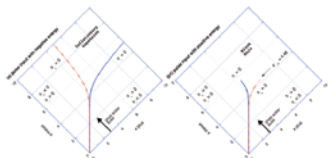


Figure 5: Horizon location, $r = r^*$, for perturbed spacetime. $\psi(r,0)$ is the case we regularize the ghost field, $\psi = 0$, and [2] and [3] are when we reduce the field, $\psi = -0.5$ and -0.1 . Dotted lines and solid lines are $r = 0$ and $r = 1$ respectively. In 4d case, the plot like the wormhole throat at $r^* = 0.5$. A complete Penrose diagram of the spacetime is in paper Torii & Shinkai.

Travel through a Wormhole – with Maintenance Operations

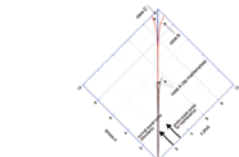


Figure 6: A total of wormhole spacetime. After a normal matter pulse, we applied a ghost wave pulse to extend the life of the wormhole throat. The timelike pulse are consistently represented with a normal scalar field pulse, $\psi_{(t,r)} = (4, 1, 1, 0, 0)$. Figure 7: A total of wormhole spacetime. After a normal matter pulse, we applied a ghost wave pulse to extend the life of the wormhole throat. The timelike pulse are consistently represented with a normal scalar field pulse, $\psi_{(t,r)} = (4, 1, 1, 0, 0)$. Figure 8: A total of wormhole spacetime. After a normal matter pulse, we applied a ghost wave pulse to extend the life of the wormhole throat. The timelike pulse are consistently represented with a normal scalar field pulse, $\psi_{(t,r)} = (4, 1, 1, 0, 0)$. Figure 9: A total of wormhole spacetime. After a normal matter pulse, we applied a ghost wave pulse to extend the life of the wormhole throat. The timelike pulse are consistently represented with a normal scalar field pulse, $\psi_{(t,r)} = (4, 1, 1, 0, 0)$.

Gauss-Bonnet gravity

- Action

$$S = \int_{\mathcal{M}} d^N x \sqrt{-g} \frac{1}{2\kappa^2} [\alpha_1 R + \alpha_2 C_{ab}C^{ab} + C_{Gauss-Bonnet}] \quad (1)$$

$$\text{where } C_{GB} = R^2 - 4R_{\mu\nu}R^{\mu\nu} + R_{\mu\nu\rho\sigma}R^{\mu\nu\rho\sigma}$$

- Field eqs.

$$\alpha_1 G_{\mu\nu} + \alpha_2 H_{\mu\nu} + g_{\mu\nu}\Lambda = \kappa^2 T_{\mu\nu} \quad (2)$$

$$\text{where } H_{\mu\nu} = 2[R_{\mu\nu} - 2R_{\mu\nu}R_{\alpha\beta} - 2R^{\alpha\beta}R_{\mu\nu\alpha\beta} + R_{\alpha\beta\gamma\delta}R^{\alpha\beta\gamma\delta}] - \frac{1}{2}g_{\mu\nu}C_{GB}$$

- Assumptions

- 5-dim. Spherical Symmetry
- Dual-null coordinate

Field Eqs.

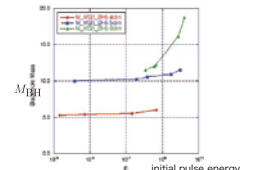
$$ds^2 = -2e^{-2\psi}dt^2 + e^{2\psi}dx^2 + r^2(d\theta^2 + \sin^2\theta d\phi^2) \quad (3)$$

- conformal factor $\Omega = \frac{1}{r}$ (4)
- expansions $\theta_{2-} = 3\partial_r r = -3\Omega^{-1}\partial_r \Omega$ (5)
- infinities $v_{2-} = \partial_t f$ (6)
- momenta of ϕ $p_{2-} = r\partial_r \phi = \Omega^{-1}\partial_r \phi$ (7)
- momenta of ψ $p_{2-} = r\partial_r \psi = \Omega^{-1}\partial_r \psi$ (8)

$$\square \text{ matter} = \text{normal field } \psi(u, r) \text{ and/or ghost field } \phi(u, r) \quad (9)$$

$$T_{\mu\nu} = T_{\mu\nu}^{\psi} + T_{\mu\nu}^{\phi} = \left[\frac{1}{2}(\partial_\mu \psi)^2 + V_1(\psi) \right] + \left[-\frac{1}{2}(\partial_\mu \phi)^2 + V_2(\phi) \right]$$

Minimum Mass of Black Hole



Quasi-Local Energy (Misner-Sharp mass in N-dim.)

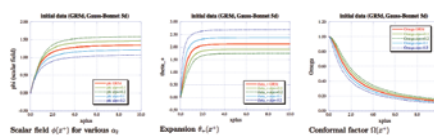
Using the area of the null 3-dimensional space of constant coordinates, \mathcal{S}_t , (or the associated area $A_t = 4\pi r_t^2$, \mathcal{E}_t may be defined as (cf Misner & Newman, 2008)

$$\mathcal{E}_t = \frac{(n-2)A_t}{2\kappa^2} \left[\frac{1}{2} \dot{\chi}^2 + \left(\chi + \frac{2}{n-2} \psi \right)^2 \right]$$

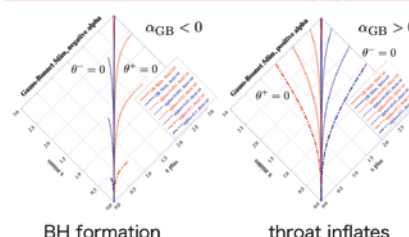
WH evolution in 5-dim. GB

in prep.

Wormholes in Gauss-Bonnet gravity (initial data on x^+)



5d GR vs Gauss-Bonnet instability appears



BH formation (throat inflates)

$$S = \int_{\mathcal{M}} d^N x \sqrt{-g} \frac{1}{2\kappa^2} [\alpha_1 R + \alpha_2 C_{ab}C^{ab} + C_{Gauss-Bonnet}] \quad (1)$$

$$\text{where } C_{GB} = R^2 - 4R_{\mu\nu}R^{\mu\nu} + R_{\mu\nu\rho\sigma}R^{\mu\nu\rho\sigma}$$